

# Accelerator Based Neutrino Experiments

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# Topics

- Neutrino Properties
  - Neutrino Flavor Oscillations
  - Mass and Mixing
  - Mass hierarchy and CP-violation
- Neutrino Oscillation Experiments using Accelerators
  - Why long baseline?
  - Why underground?
  - Why so large?

# Neutrinos are abundant in our universe

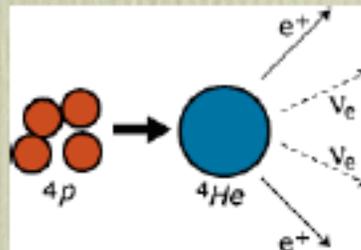
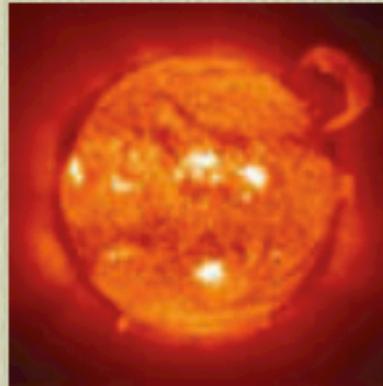


*Remnants of the Big Bang  
(~1000 per cubic cm)*

*Supernovae  
(99% of energy is emitted  
as neutrinos)*

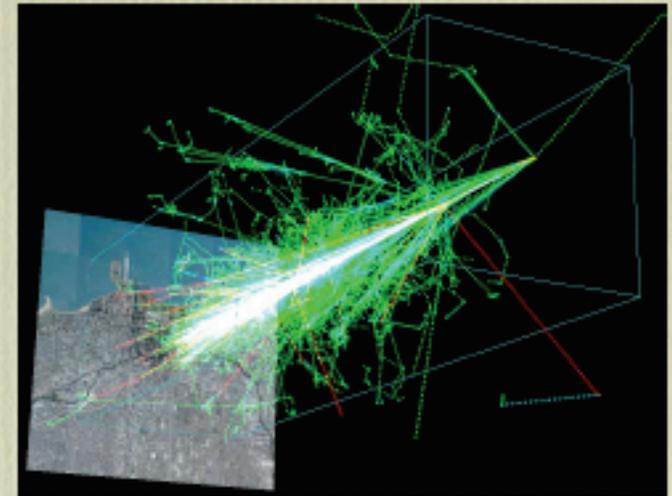


*Our Sun is powered by  
nuclear reactions*



*(100 billion per sq cm every  
second)*

**Cosmic ray interactions**



**Nuclear Power Plants**



... but difficult to detect



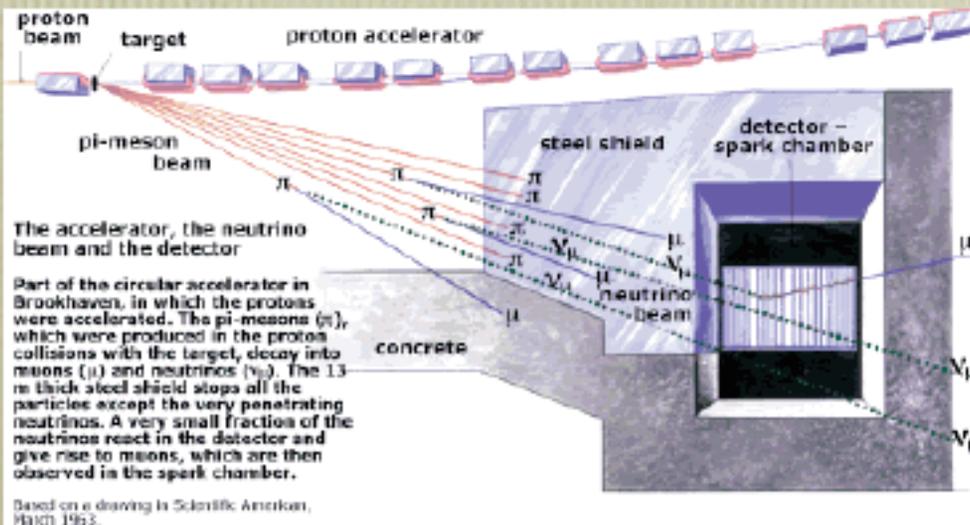
postulated in 1932



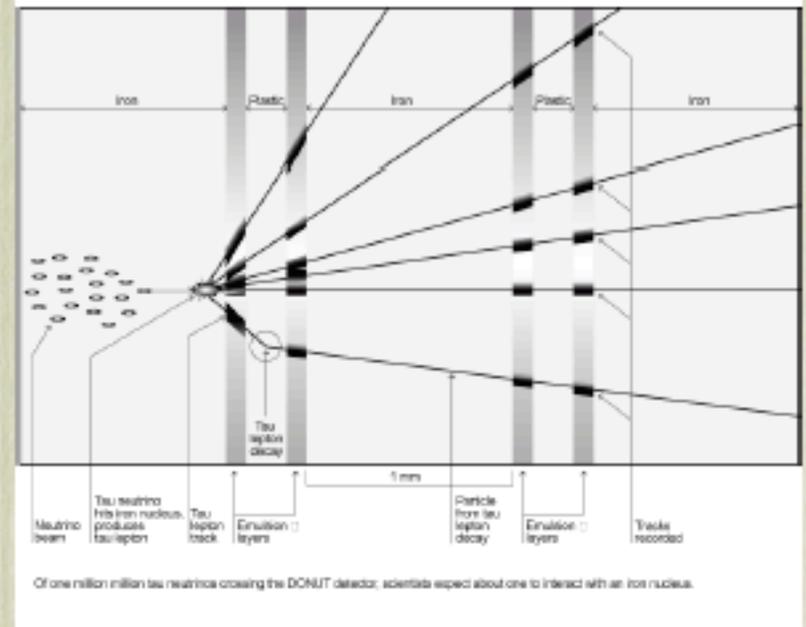
1st detected in 1956

**IN 2000  
A GROUP OF  
PHYSICISTS FINALLY  
FOUND EVIDENCE OF  
THE TAU TYPE OF  
THIS SUBATOMIC  
PARTICLE**

2 neutrino hypothesis (-1947) confirmed in 1961



### Detecting a Tau Neutrino



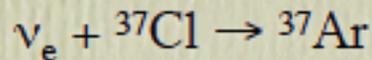
# The Solar Neutrino Mystery



100,000 gallons of perchlorethylene  
 $10^{30}$  atoms of chlorine

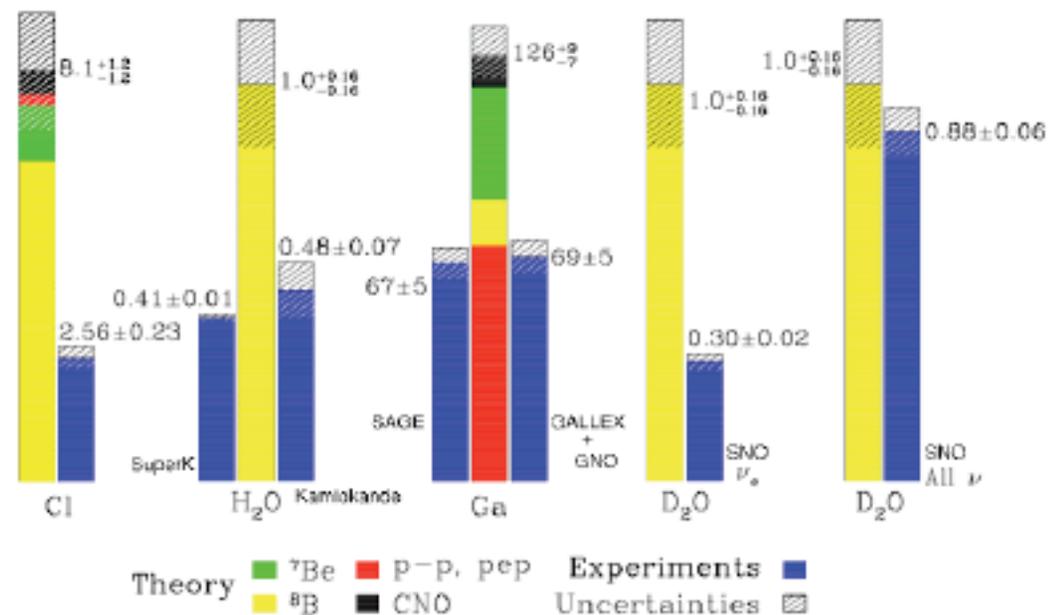
Over the next 30 years the deficit  
 would be observed in many  
 different detectors

~1970 : detecting solar neutrinos via the inverse  
 $\beta$ -capture

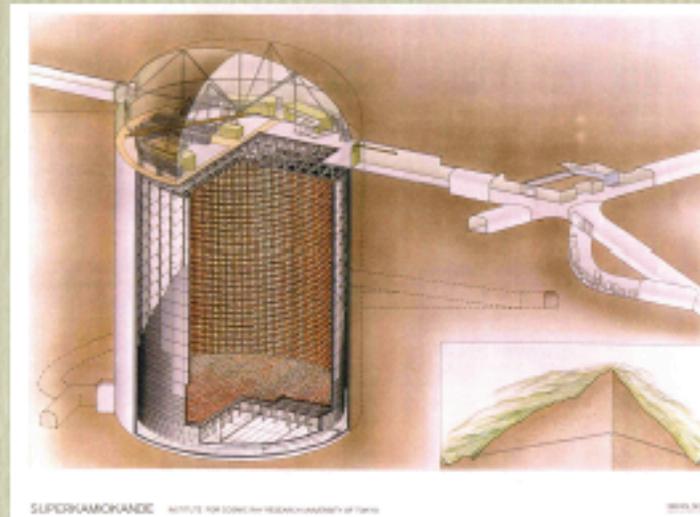
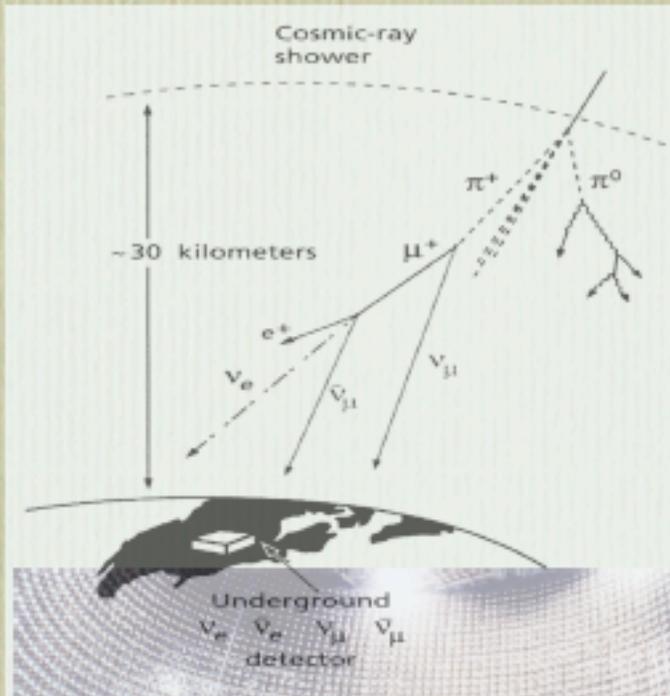


Only about 1/3 of the number expected were  
 observed.

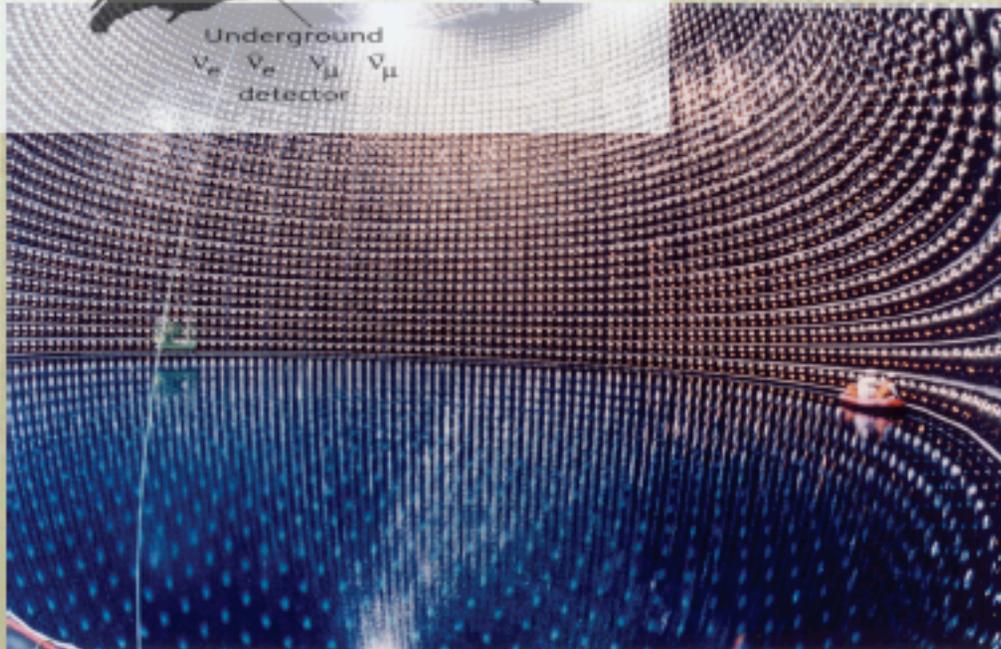
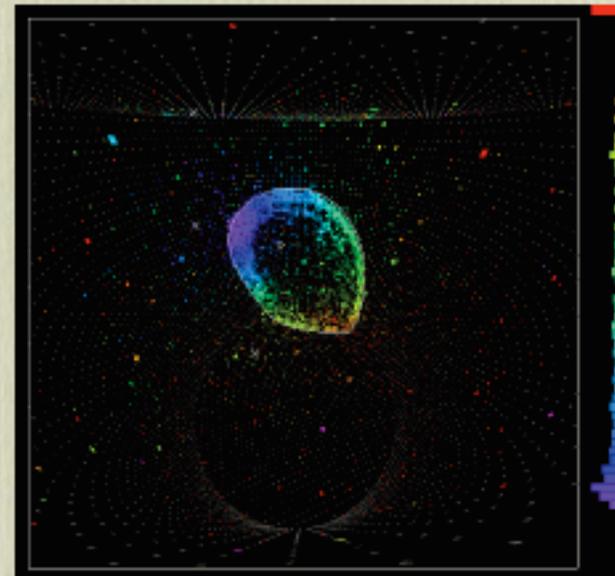
Total Rates: Standard Model vs. Experiment  
 Bahcall-Serenelli 2005 [BS05(OP)]



# more "missing" neutrinos .....



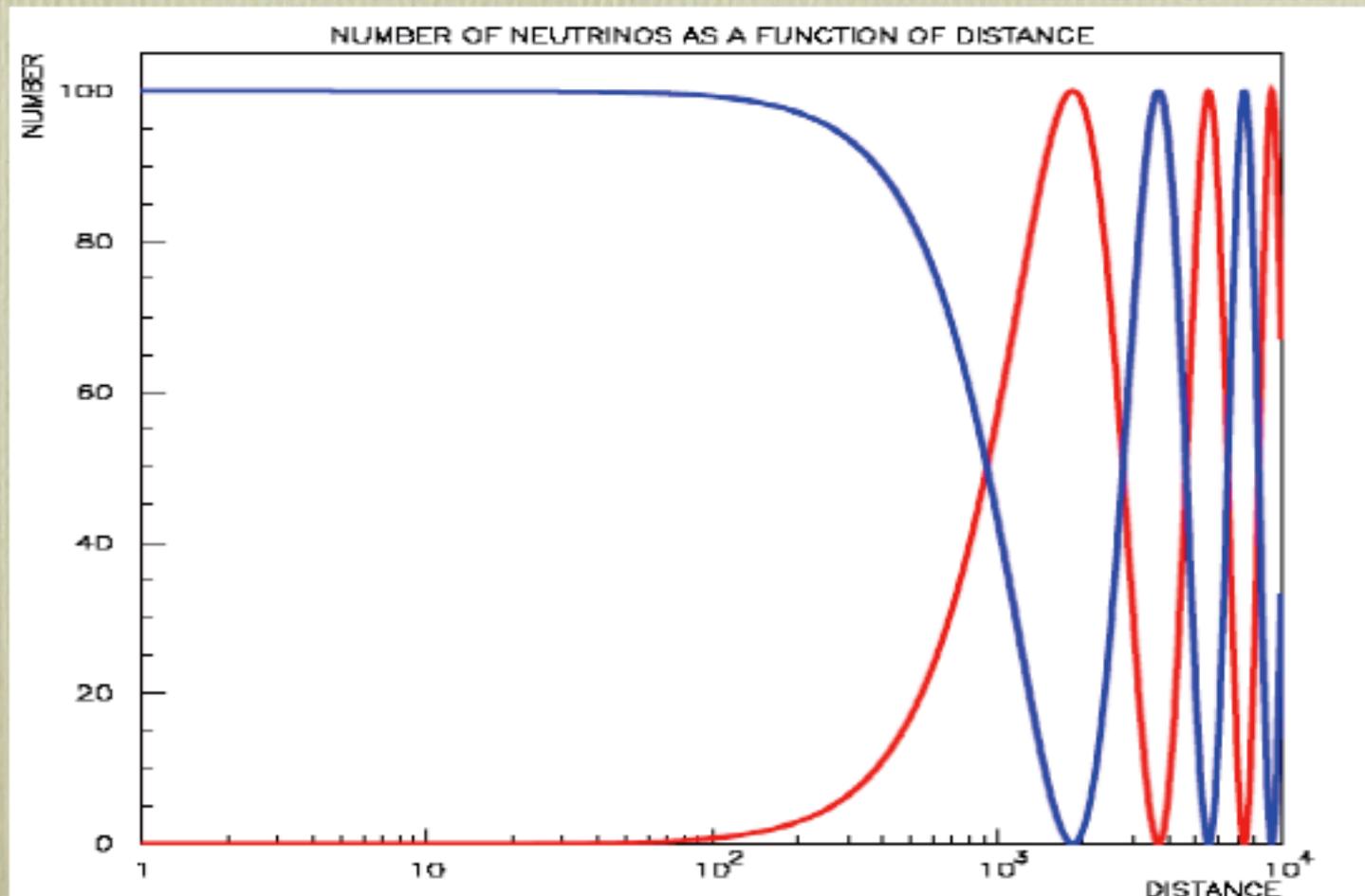
**1/2 the predicted # of muon neutrinos are not detected**



# The Hypothesis : neutrinos have mass; oscillate between mass states

$$P(\nu_a \rightarrow \nu_b) = \alpha \cdot \sin^2\left(\frac{1.267 \cdot \Delta m^2 \cdot L}{E}\right)$$

Oscillation Probability depends on :  
neutrino energy  
mass difference  
distance traveled



*Solar and atmospheric data indicated that if the missing neutrinos were indeed due to oscillations, then **the mass differences were very small** :*

$$\Delta m_{12}^2 \approx 8 \times 10^{-5} \quad \Delta m_{23}^2 \approx 3 \times 10^{-3}$$

*Knowing these mass scales guides the design of a controlled laboratory experiment.*

# Neutrinos can be made at a proton accelerator : i.e. Fermilab Main Injector

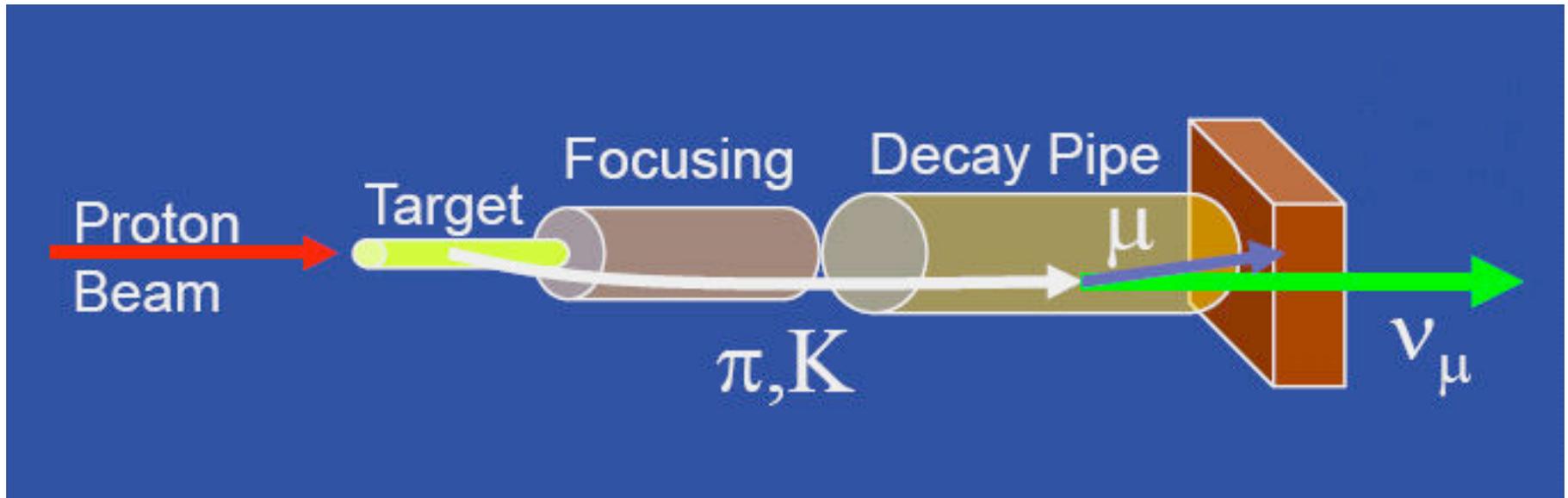


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$\sim 10^{14}$  protons, travelling near the speed of light, can be extracted every 1 – 2 seconds

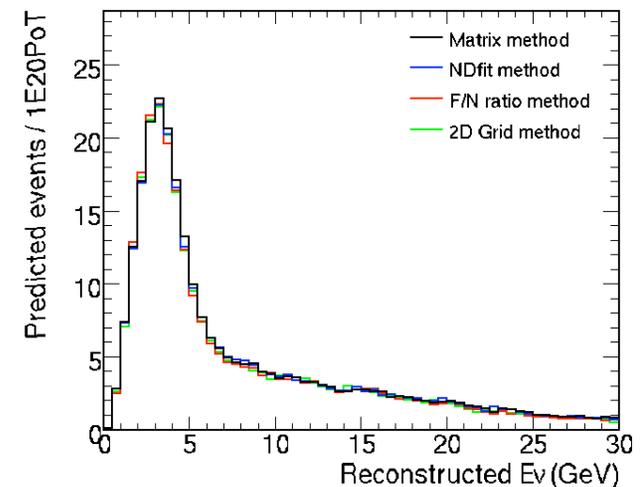
# Ingredients in a neutrino beam



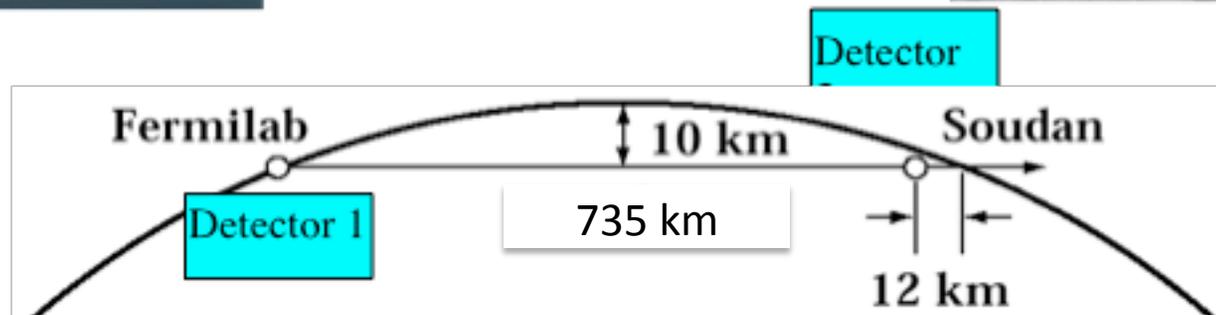
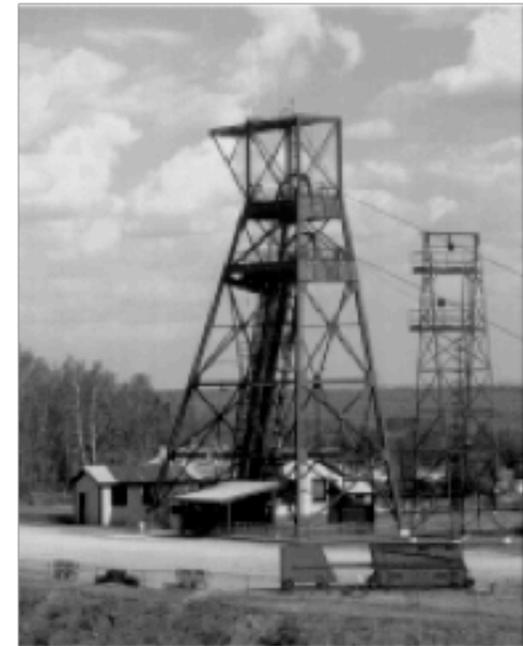
At Fermilab we can produce  
 $\sim 10^{11}$  neutrinos/m<sup>2</sup>/10<sup>14</sup> protons

$$P(\nu_\mu \rightarrow \nu_\tau) = f(\Delta m_{23}^2, E_\nu, L)$$

A blue bracket is drawn under the  $\Delta m_{23}^2$  term in the equation, with a blue arrow pointing from the bracket towards the plot below.

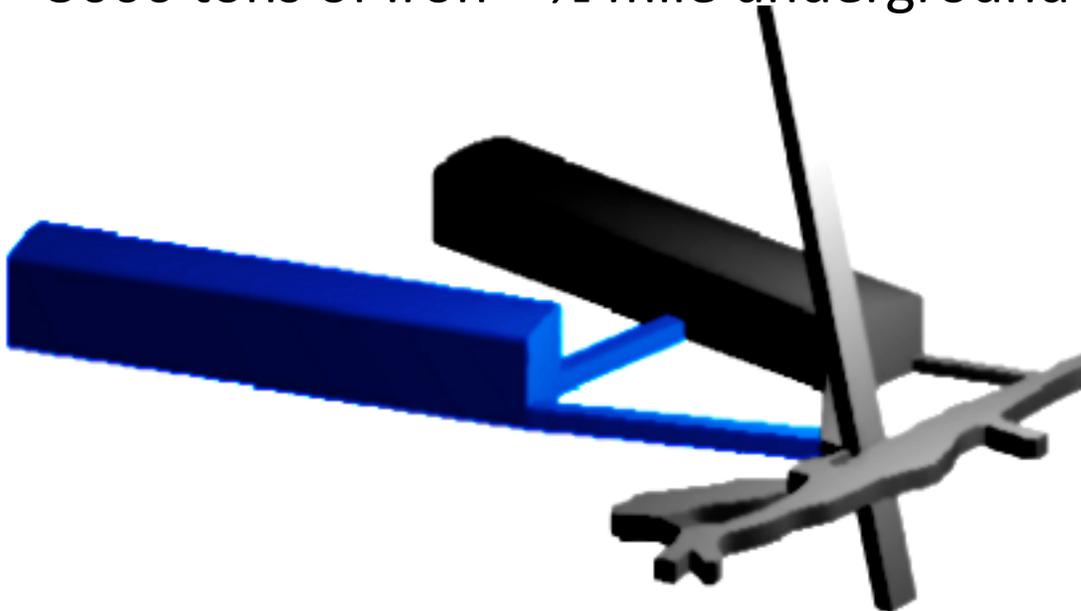


# MINOS - Main Injector Neutrino Oscillation Search (Minnesota-Illinois)



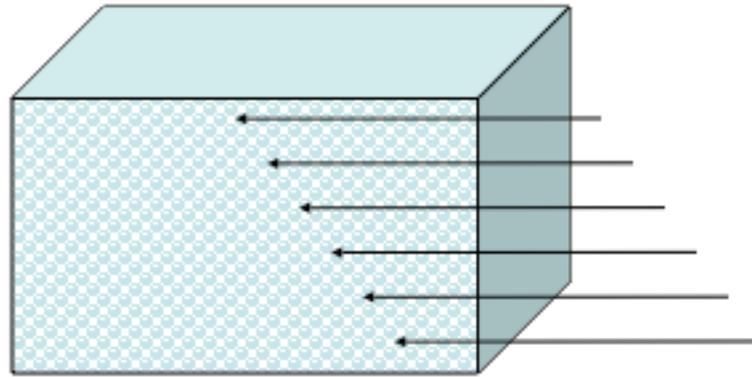


5000 tons of Iron – ½ mile underground



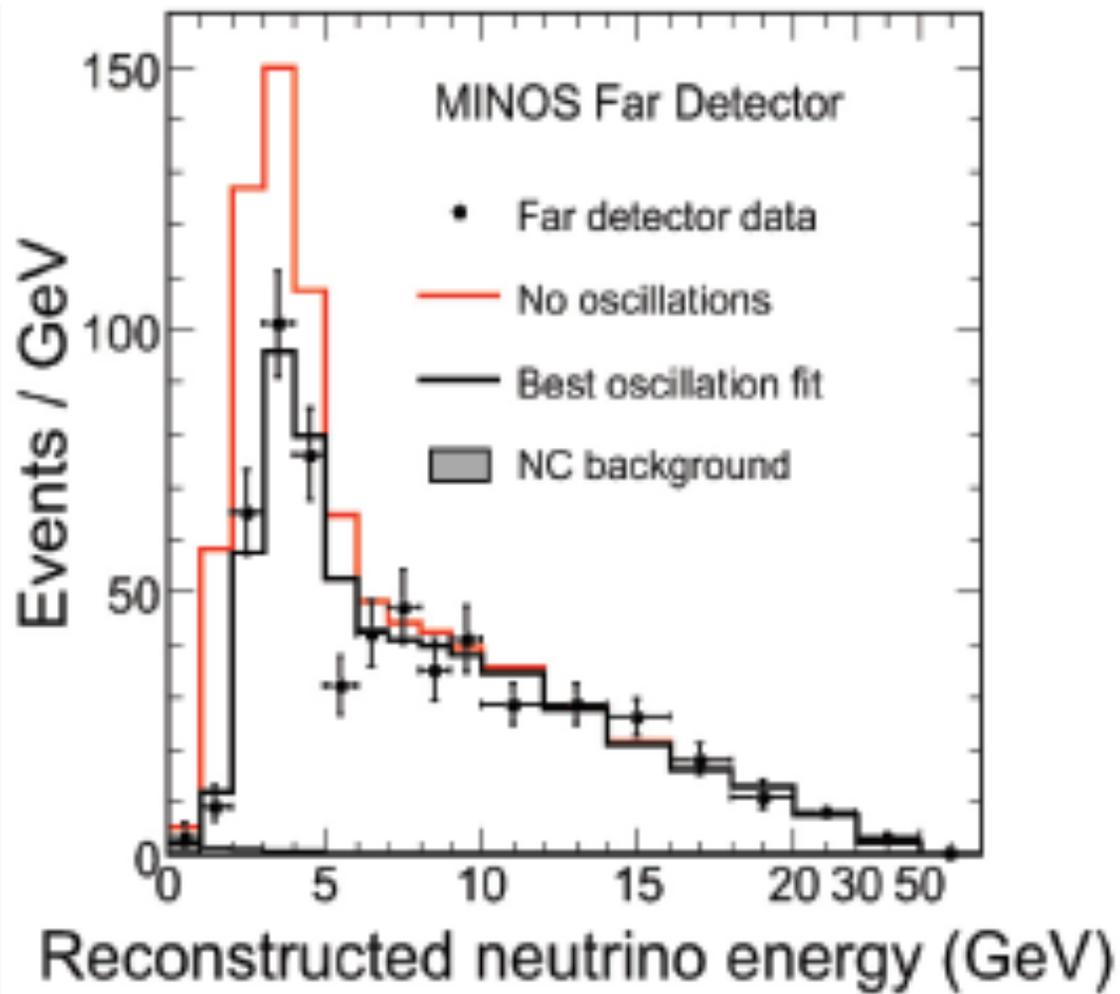
# Why is it so difficult to detect neutrinos?

Consider a “beam” of particles, incident on a mass of material (i.e. atomic scattering centers)

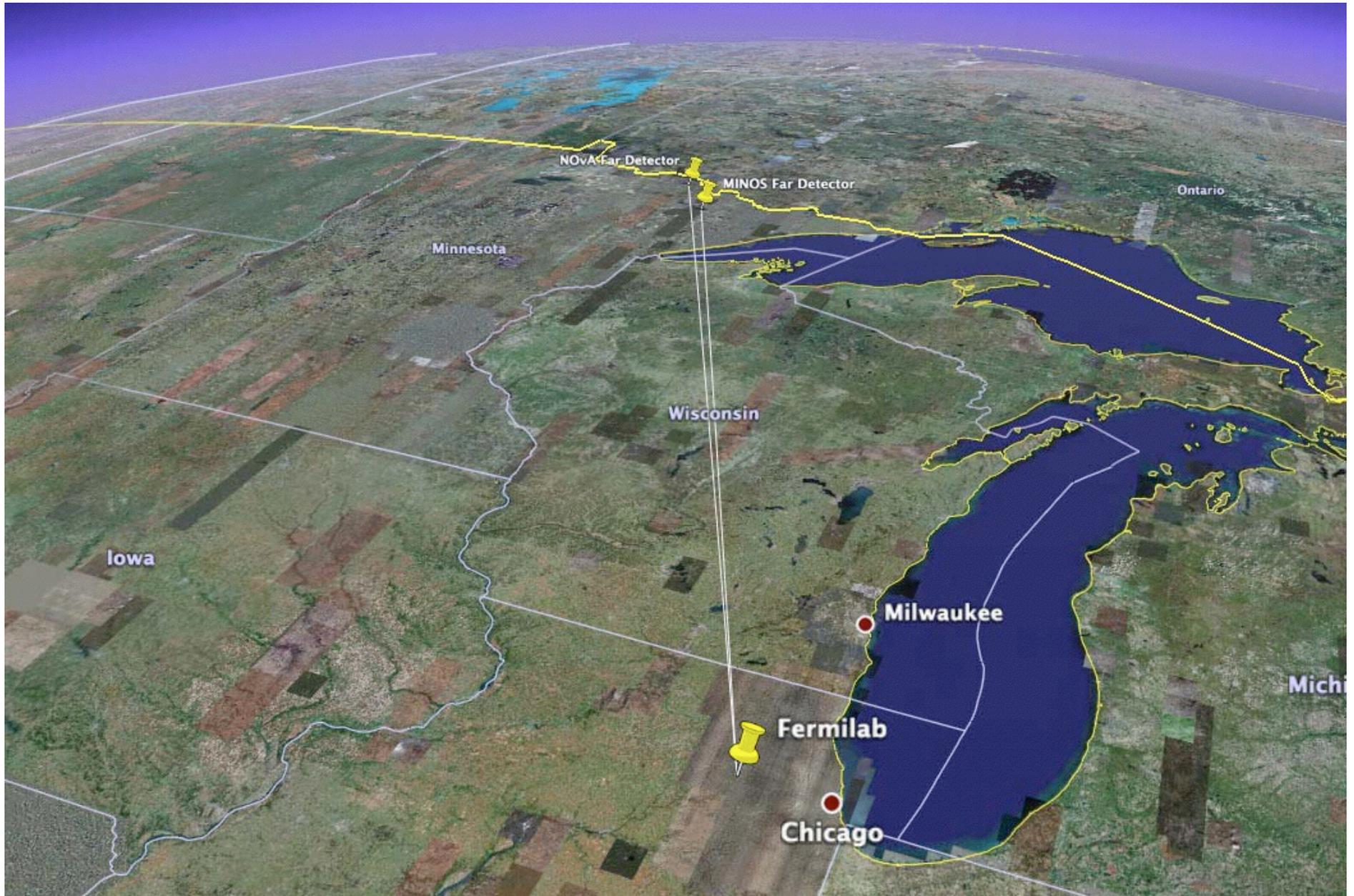


The **probability** that one of the particles will interact depends on the number of scattering centers, i.e. **protons/neutrons**, **density** of the scattering centers, the **intensity** of the beam, and the **energy** of the beam

No oscillations – expect  $\sim 3$  events/day; observe  $\sim 1-2$

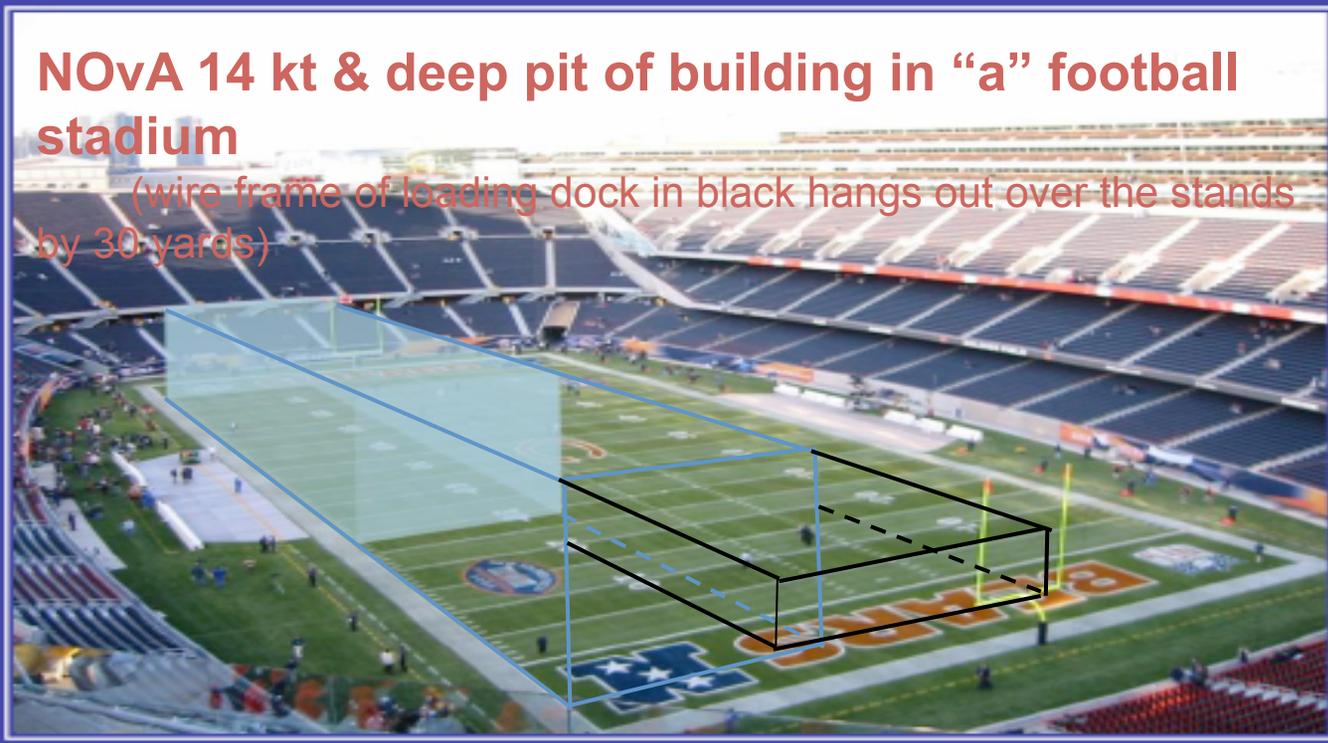


# NOvA : NuMI Off-Axis



## NOvA 14 kt & deep pit of building in “a” football stadium

(wire frame of loading dock in black hangs out over the stands by 30 yards)



Start data taking with full detector in 2013 with  $6 \times 10^{20}$  POT/year

Run 3 years producing neutrinos and 3 years producing anti-neutrinos

All goes well – significant data taking complete in 2019-2020

# Fermilab to Homestake DUSEL (1290km)



# DUSEL Deep Underground Science and Engineering Laboratory at Homestake, SD



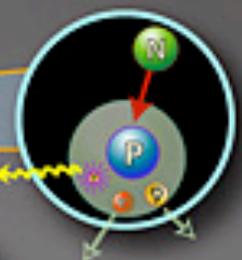
Engineering



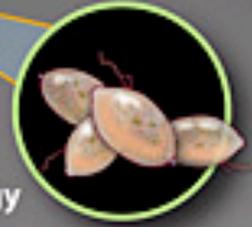
Geoscience



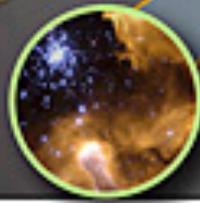
Physics



Biology



Astrophysics



6 1/2 Empire State Buildings for scale

Shallow Lab

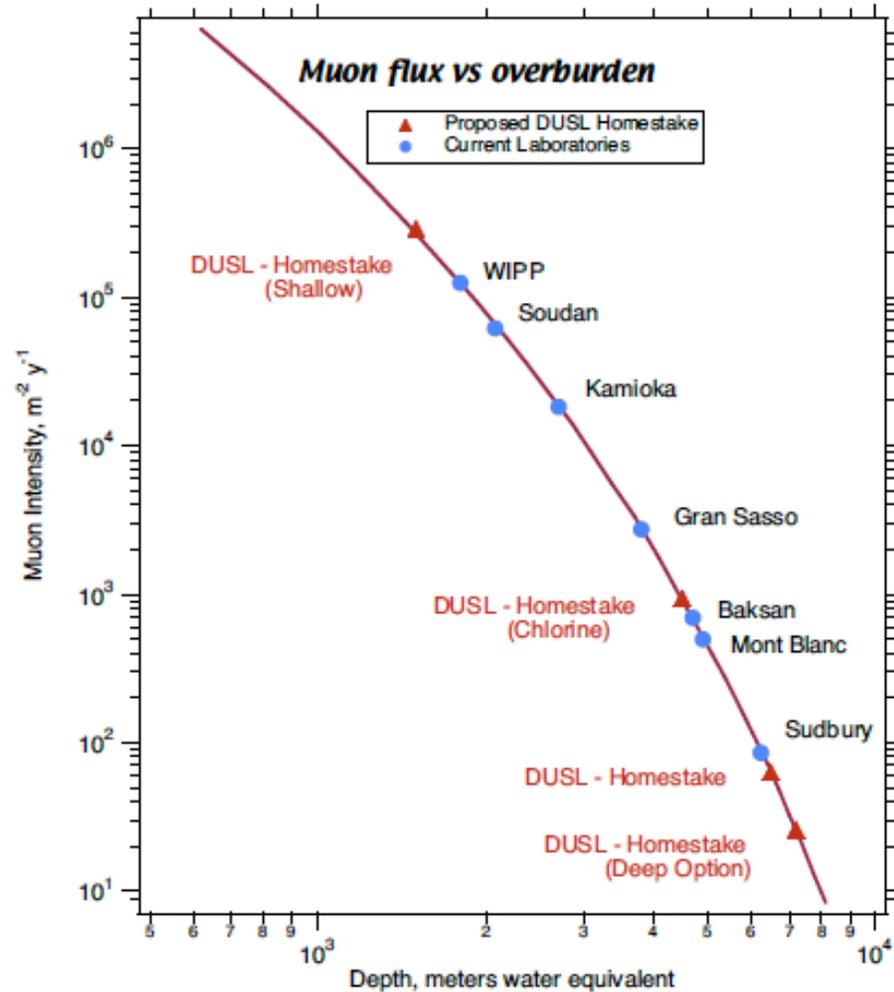
Mid-level

Deep Campus

Open cut



# Why Go Underground?



Depth reduces the rate of cosmic ray muons and associated interactions occurring in the detector which can mimic or confuse the information about the real events of interest.



Waste Water Treatment

Open Cut

Yates Complex

Town of Lead

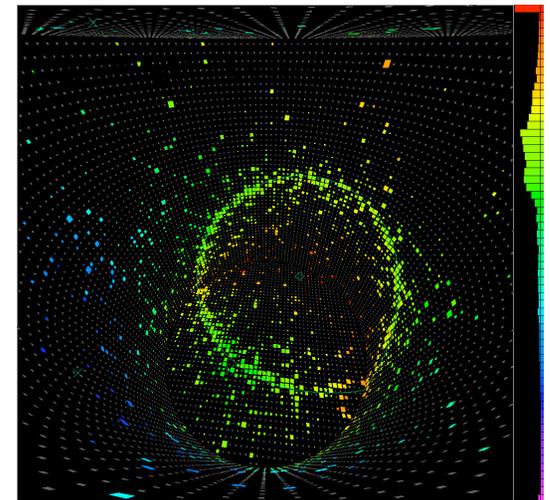
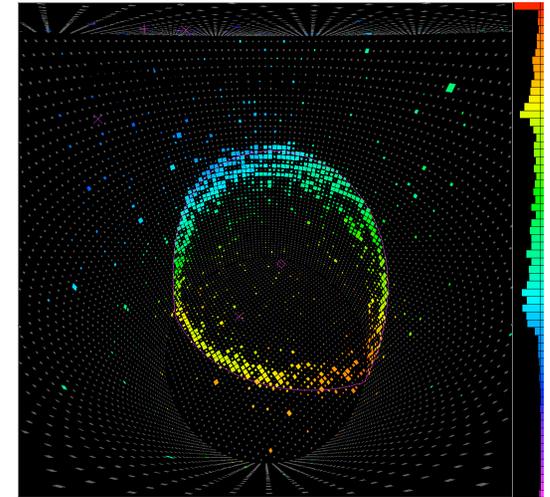
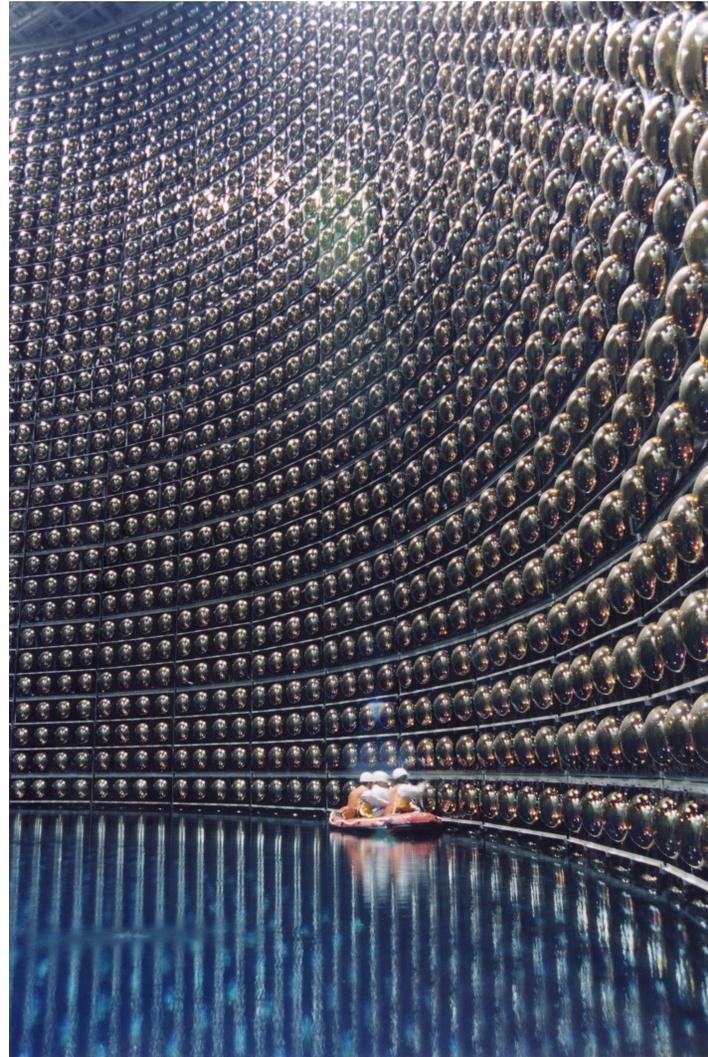
1 km

Ross Complex

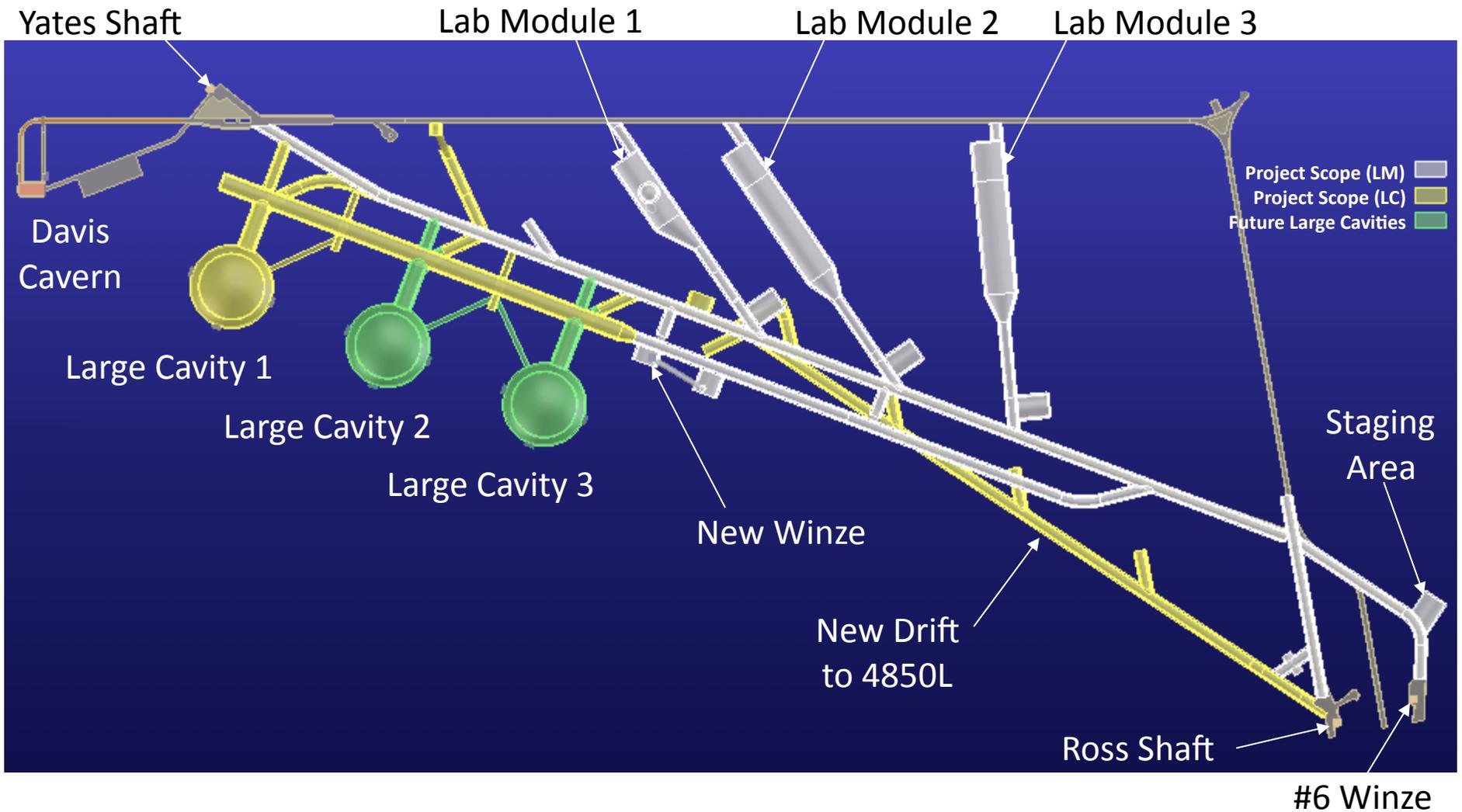
Kirk Canyon Adit

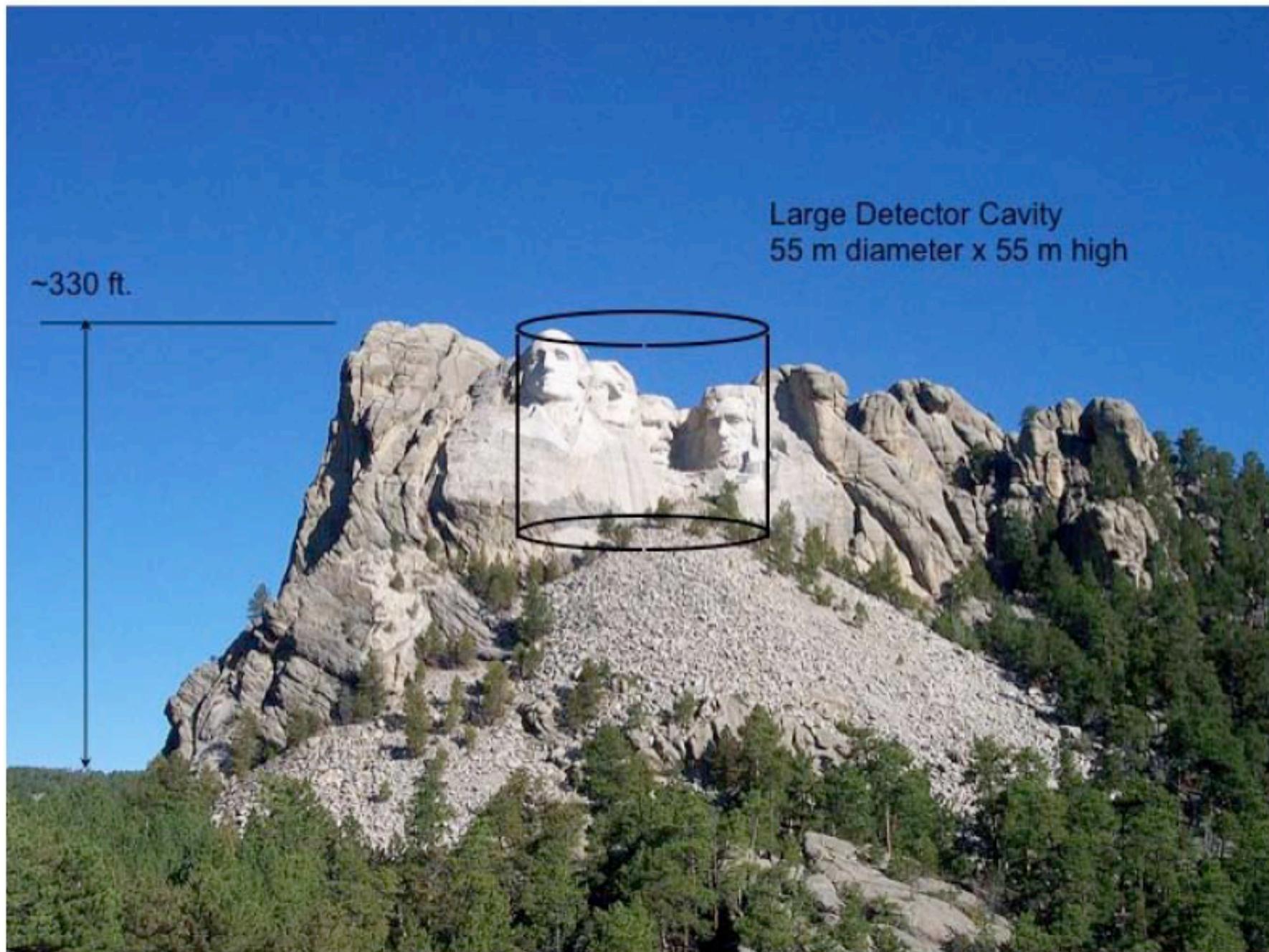
# Planning the world's largest detectors

- Super-K (taking data in Japan)
  - 13K 20" PMT
  - 40% coverage
  - 50 kT total mass
  - 39 m diameter
  - 42 m height
- LBNE - proposed
  - 60 K 10" PMT per 100kT FV module (25%)
  - ~55 m diameter
  - ~60 m height



# 4850 Level Developmental Baseline for PDR: Three Lab Modules & Three Large Cavities, Plan View





~330 ft.

Large Detector Cavity  
55 m diameter x 55 m high



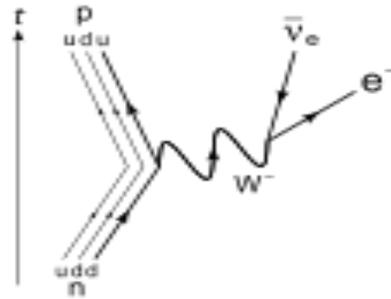
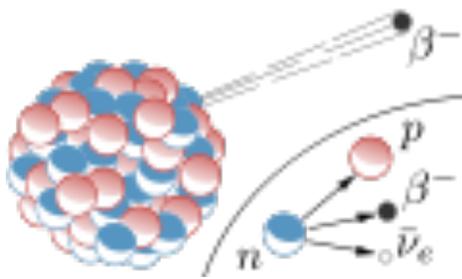
# Summary

- Experiments to study nature's tiniest particles require tools of unprecedented scale.
- By producing neutrinos at an accelerator we are able to systematically study the properties of the neutrinos, in particular the behavior of the neutrino versus anti-neutrino.
- These experiments might possibly lead us to uncover one of the most compelling questions about our universe.

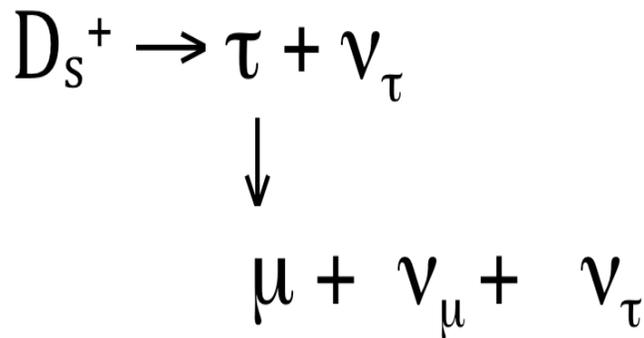
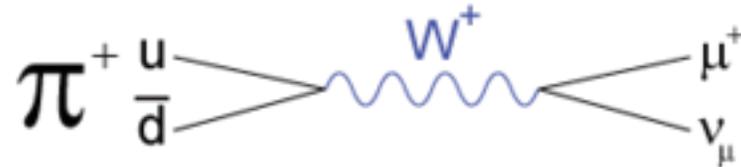
*Thank you*

# Backup Material

# Neutrino Production



Radioactive b-decay



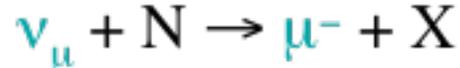
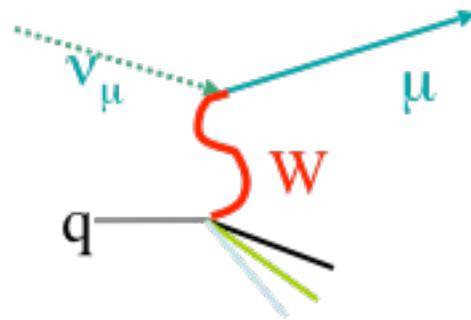
*Neutrinos are produced in the weak interaction  
Typically from the decay of radioactive nuclei,  
pion and other sub-atomic particles, from both  
natural and man made sources...*

# Neutrino Flavors & Interactions

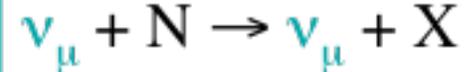
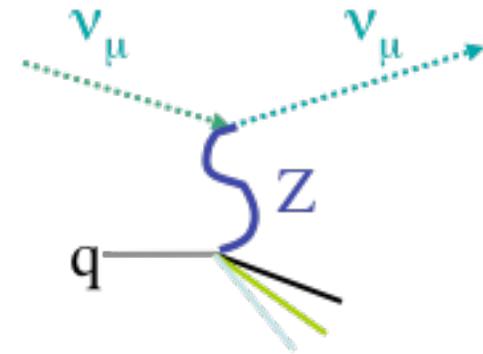
Elementary Particles				
Leptons	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	Force Carriers
	$e$ electron	$\mu$ muon	$\tau$ tau	
Quarks	$u$ up	$c$ charm	$t$ top	Force Carriers
	$d$ down	$s$ strange	$b$ bottom	
				$\gamma$ photon
				$g$ gluon

I    II    III

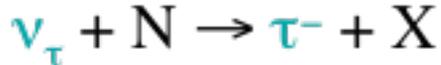
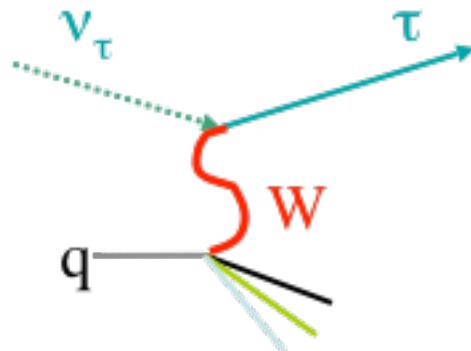
Three Families of Matter



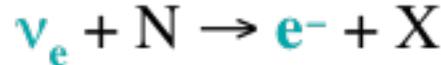
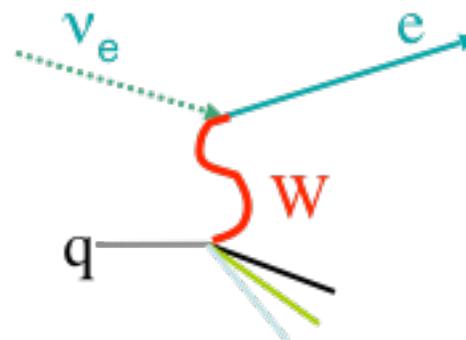
*Charged current*



*Neutral current*

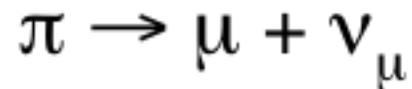


*Tau Charged current*



*Electron Charged current*

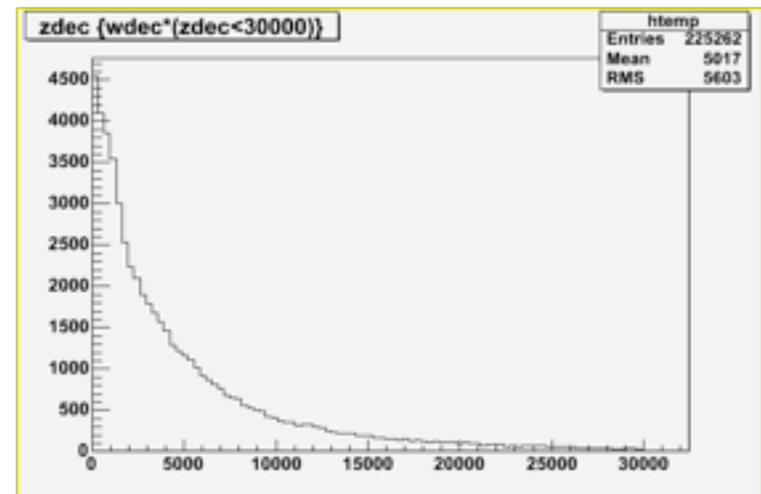
# Neutrinos come from $\pi$ decay



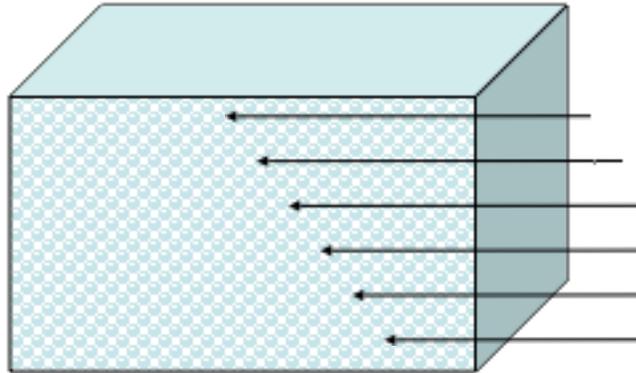
$$\tau_{\pi} = 2.6 \times 10^{-8} \text{ sec}$$

The average distance,  $d$ , traveled by an unstable, relativistic particle before it decays is given by  $d = \gamma c \tau$ , where  $\gamma = \frac{1}{\sqrt{1 - \left(\frac{v}{c}\right)^2}}$  and  $c$  is the speed of light

For a 10 GeV  $\pi$ ,  
 $\gamma \sim 27$ , and  $d \sim 220$  m



For a beam of neutrinos .....



$$\# \text{ of } \nu \text{ interactions} = \Phi \times \sigma/n \times N_n$$

$\Phi$  = neutrinos/unit area

$\sigma/n$  = cross section per nucleon  
(probability of interaction)

$N_n$  = number of nuclei in the target

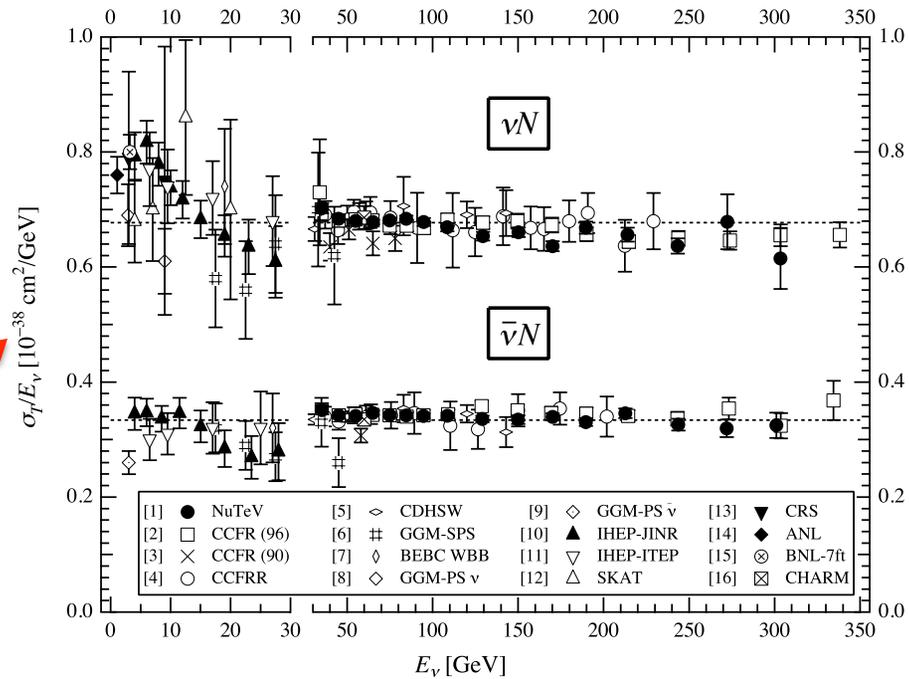
$\sim 6 \times 10^9$  nuclei in a 1000 tons of iron

Recall our flux is

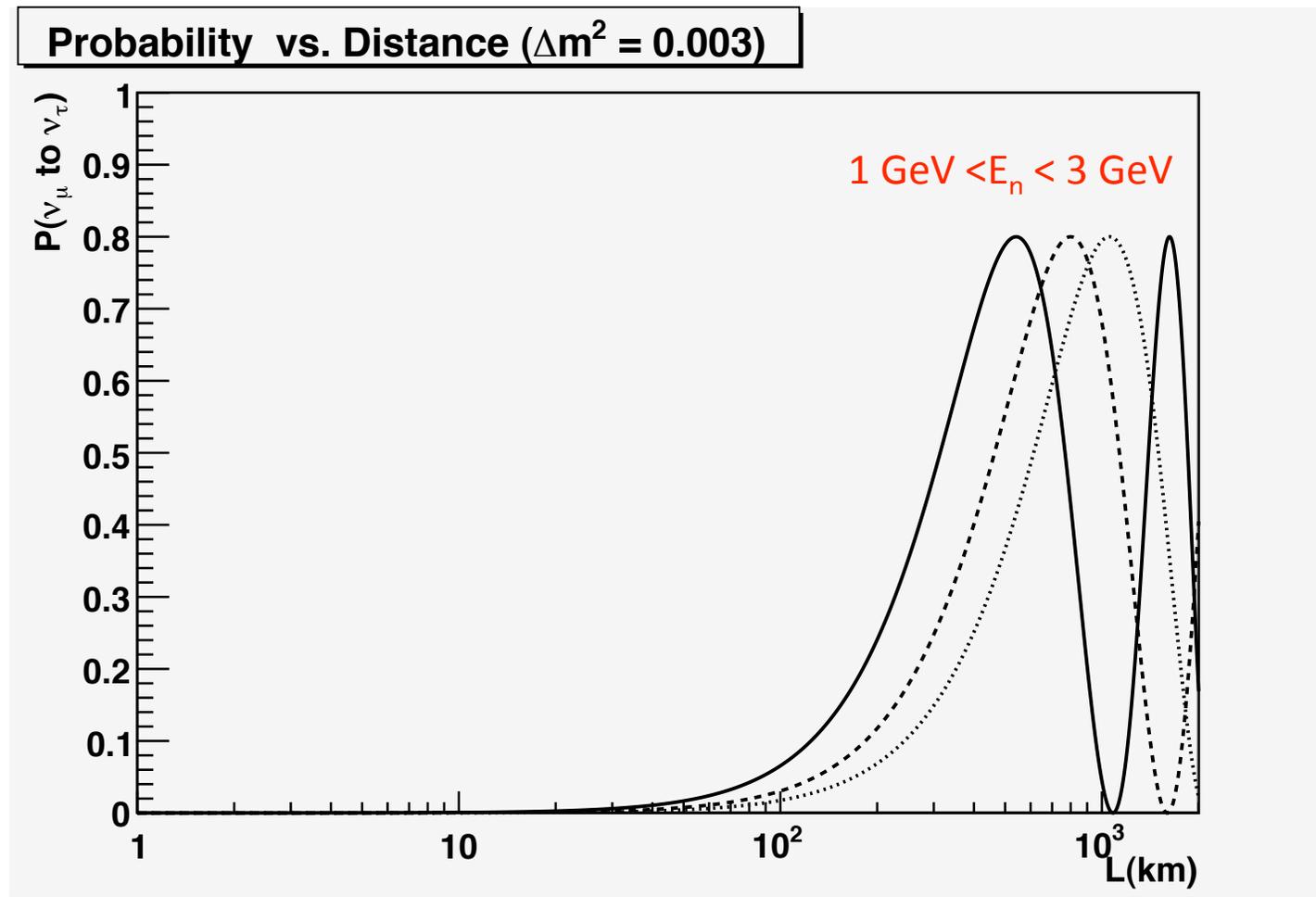
$\sim 10^{11}$  neutrinos/m<sup>2</sup>/10<sup>14</sup> protons

$10^{-38}$  ↗

It takes a lot of neutrinos (protons),  
and target material to compensate  
for the **very small cross section**



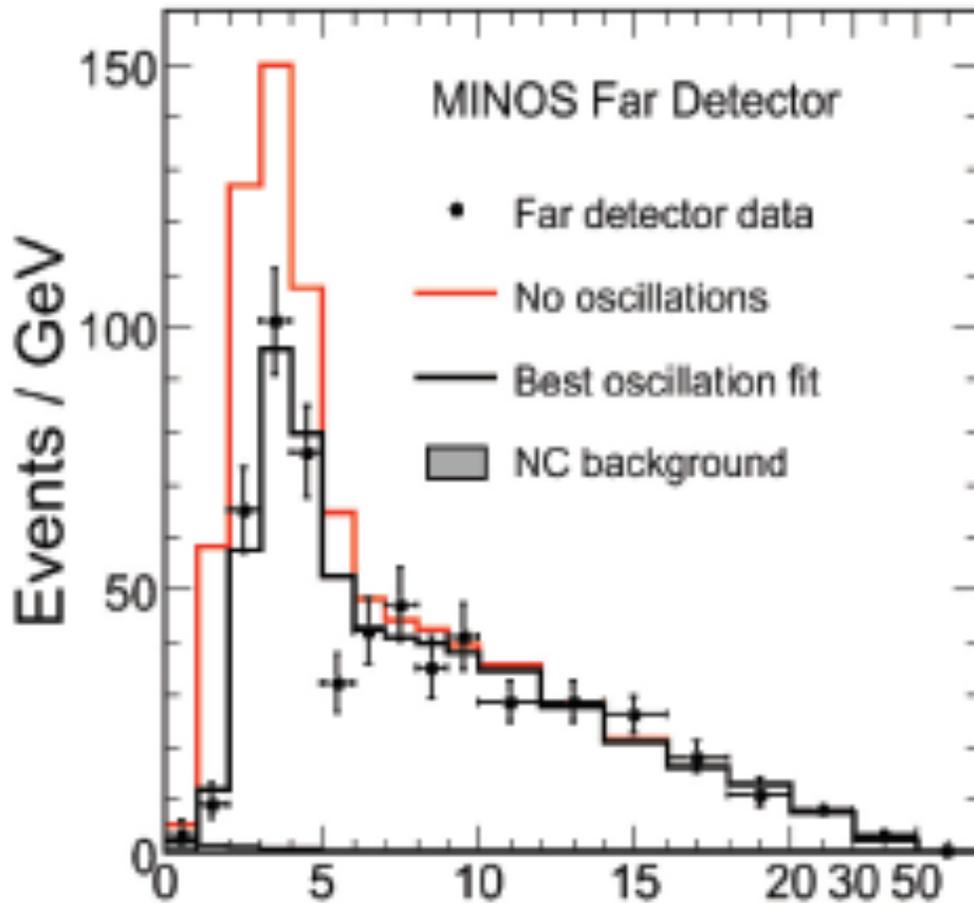
# Why do we go over a “long baseline”?



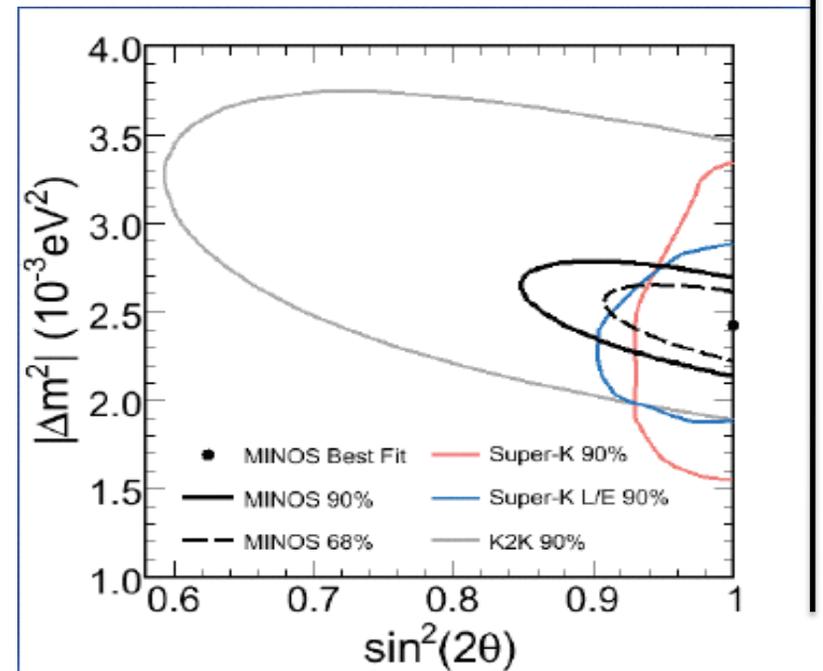
$$P(\nu_\mu \rightarrow \nu_\tau) = \sin^2(2\theta_{23})\sin^2(1.27\Delta m^2_{23} L/E)$$

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2(2\theta_{23})\sin^2(1.27\Delta m^2_{23} L/E)$$

No oscillations – expect  $\sim 3$  events/day; observe  $\sim 1-2$



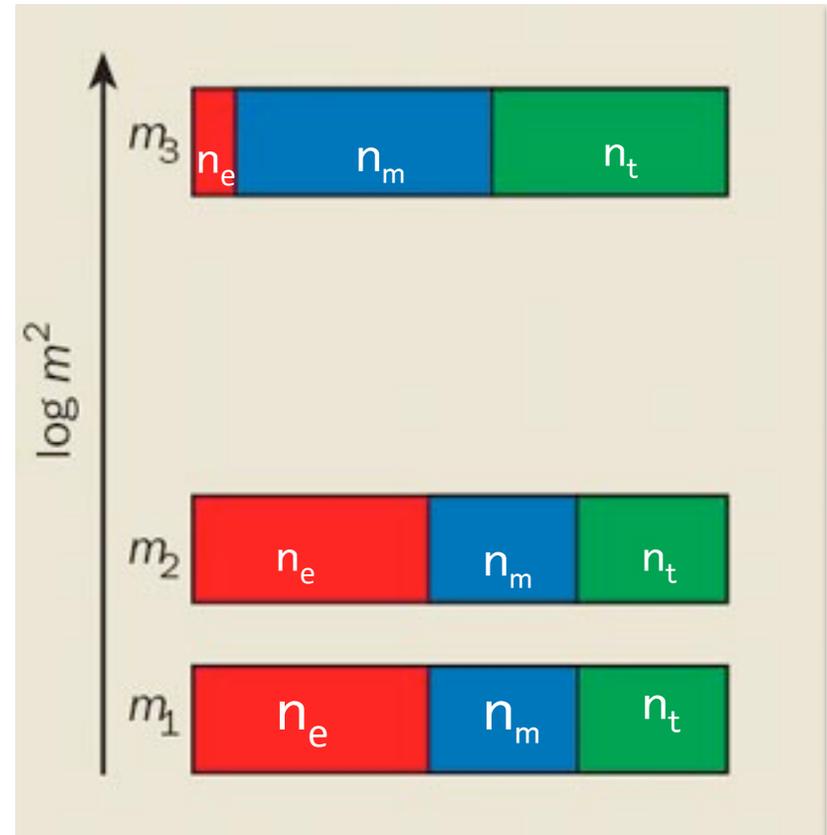
Reconstructed neutrino energy (GeV)



# Three neutrino mass and mixing

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Solar ns  $\nu_e$  appearance  
 $\nu_\mu$  disappearance (i.e. MINOS)



$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

To make matters more complicated, we don't know the ordering of the neutrino mass states



This is known as the “**neutrino mass hierarchy**” question

# Consider the Probability of

$$\nu_{\mu} \rightarrow \nu_e$$

*For neutrinos travelling in a vacuum :*

$$P(\nu_{\mu} \rightarrow \nu_e)_{vacuum} = f(\theta_{12}, \theta_{13}, \theta_{23}, \delta_{CP}, \Delta m_{12}^2, (sign)\Delta m_{23}^2, E_{\nu}, L)$$

*If CP is conserved :*

$$P(\nu_{\mu} \rightarrow \nu_e) = P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e)$$

*For neutrinos travelling through matter :*

$$P(\nu_{\mu} \rightarrow \nu_e)_{matter} = f(\theta_{12}, \theta_{13}, \theta_{23}, \delta_{CP}, \Delta m_{12}^2, (sign)\Delta m_{23}^2, E_{\nu}, L, V_{matter})$$

$$V_{matter} = f(\text{electron number density})$$

# Why does it *matter*?

$$P(\nu_\mu \rightarrow \nu_e)_{\text{matter}} \neq P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)_{\text{matter}} \\ \sim \text{fake CP violation}$$

*If  $m_3 > m_2$  (Normal mass hierarchy)*

$$P(\nu_\mu \rightarrow \nu_e)_{\text{matter}} > P(\nu_\mu \rightarrow \nu_e)_{\text{vacuum}} \\ P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)_{\text{matter}} < P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)_{\text{vacuum}}$$

*Neutrinos are enhanced and anti-neutrinos are suppressed*

*If  $m_3 < m_2$  (Inverted mass hierarchy)*

$$P(\nu_\mu \rightarrow \nu_e)_{\text{matter}} < P(\nu_\mu \rightarrow \nu_e)_{\text{vacuum}} \\ P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)_{\text{matter}} > P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)_{\text{vacuum}}$$

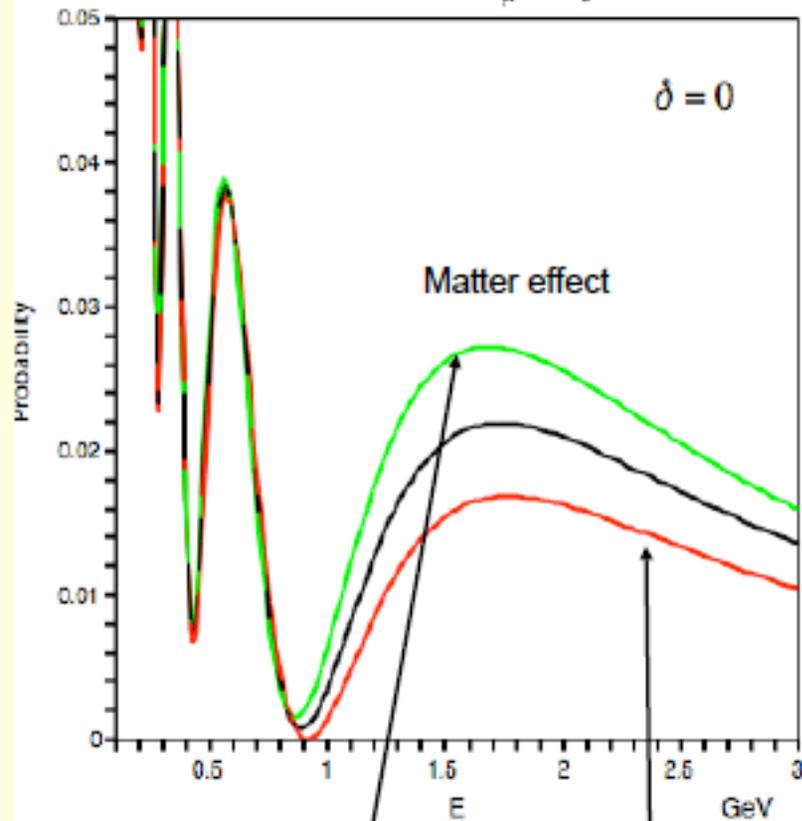
*Anti-neutrinos are enhanced and neutrinos are suppressed*

We can design experiments to try and tell these different situations apart!

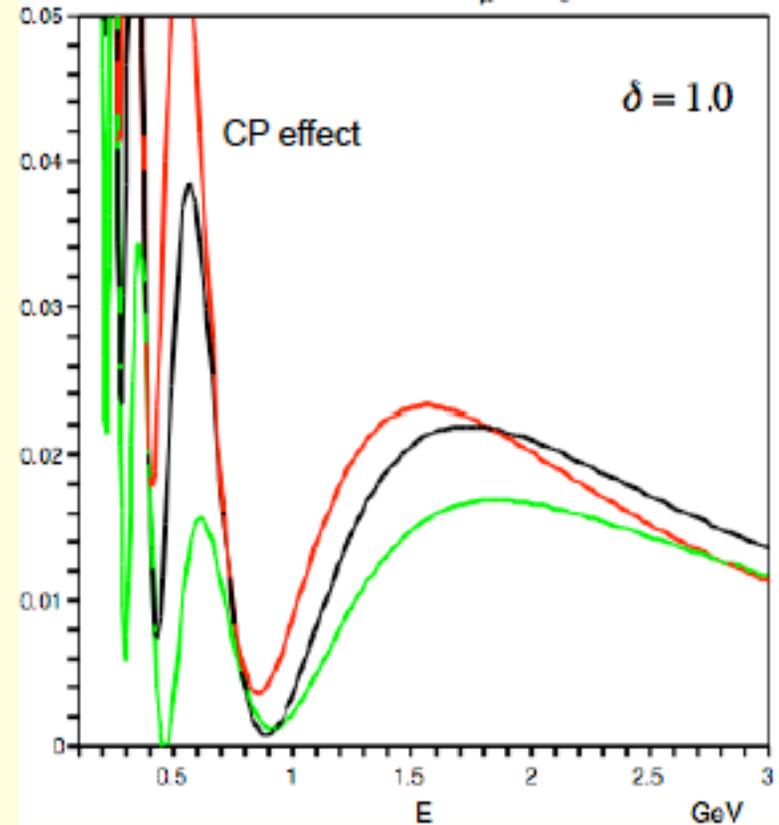
# Matter Effects and CP violation

Normal hierarchy  
 $\sin^2(2\theta_{13}) = 0.04$

810 km :  $\nu_\mu \rightarrow \nu_e$



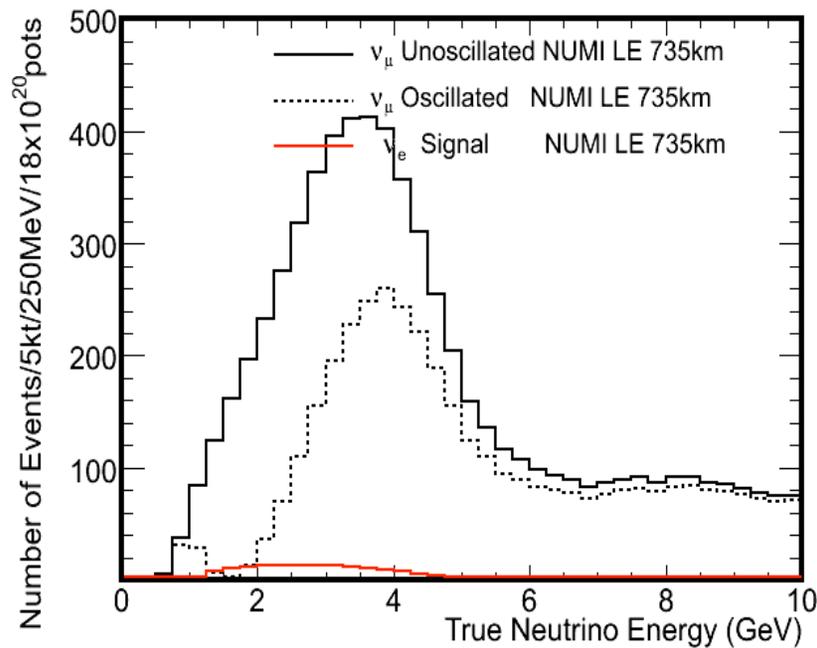
810 km :  $\nu_\mu \rightarrow \nu_e$



$\nu$ 's and  $\bar{\nu}$ 's have different oscillation probabilities

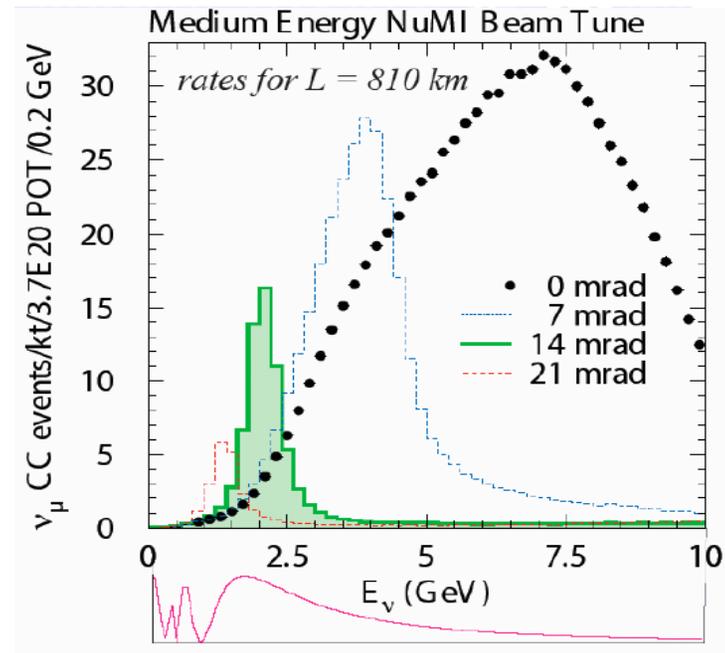
# Optimizing an experimental configuration

MINOS



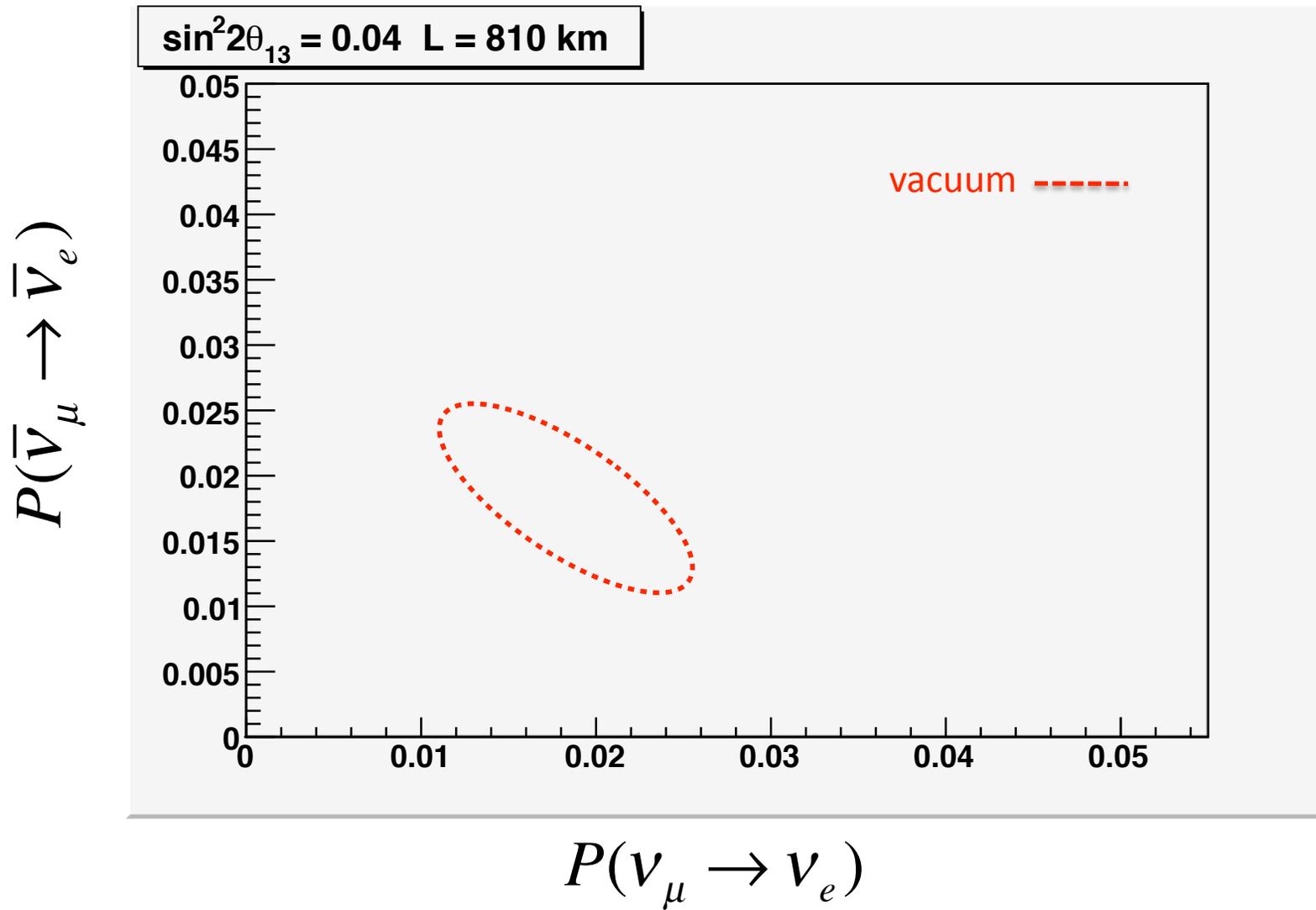
Must be able to detect muon neutrinos

NOvA

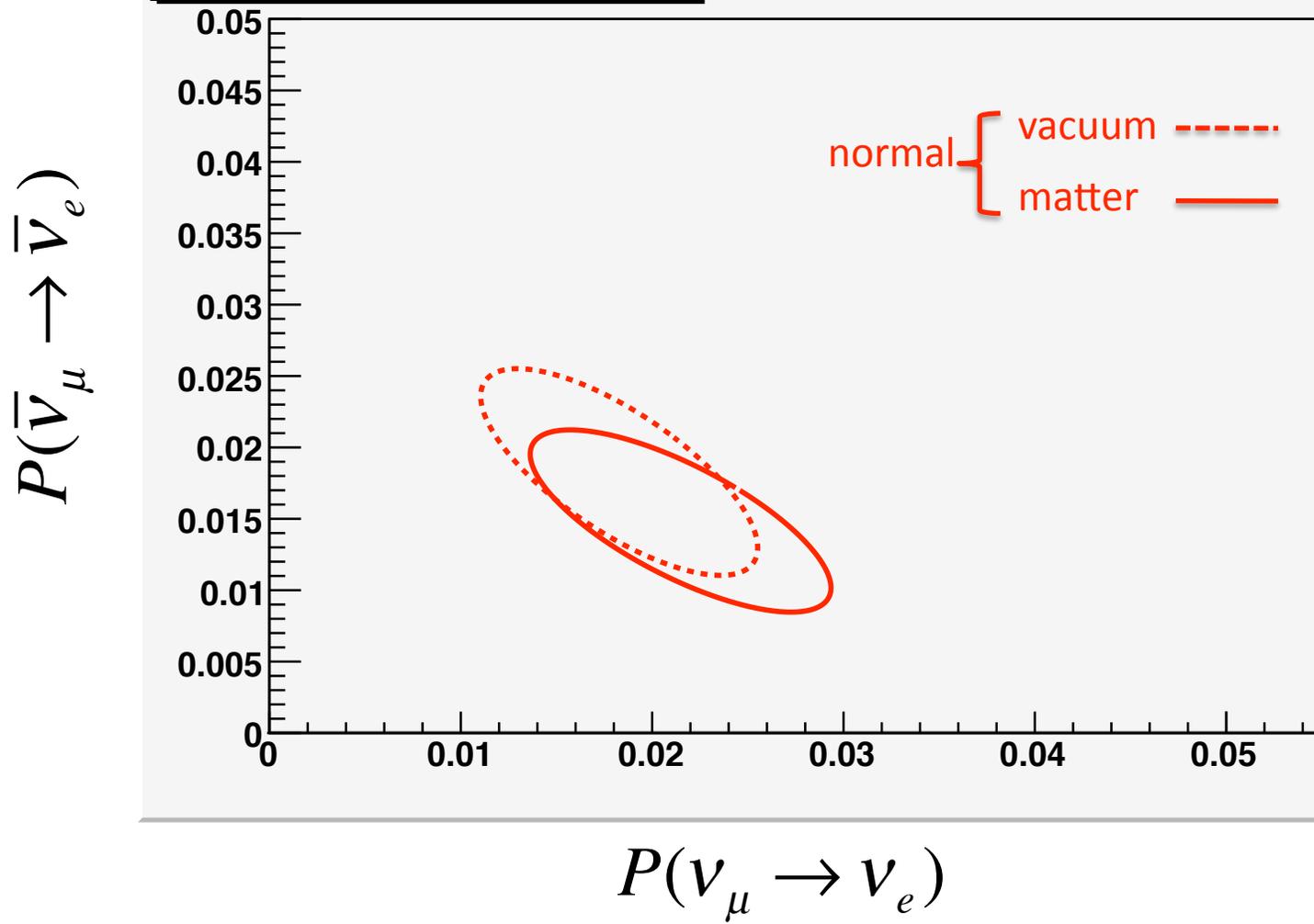


Must be able to detect electron neutrinos

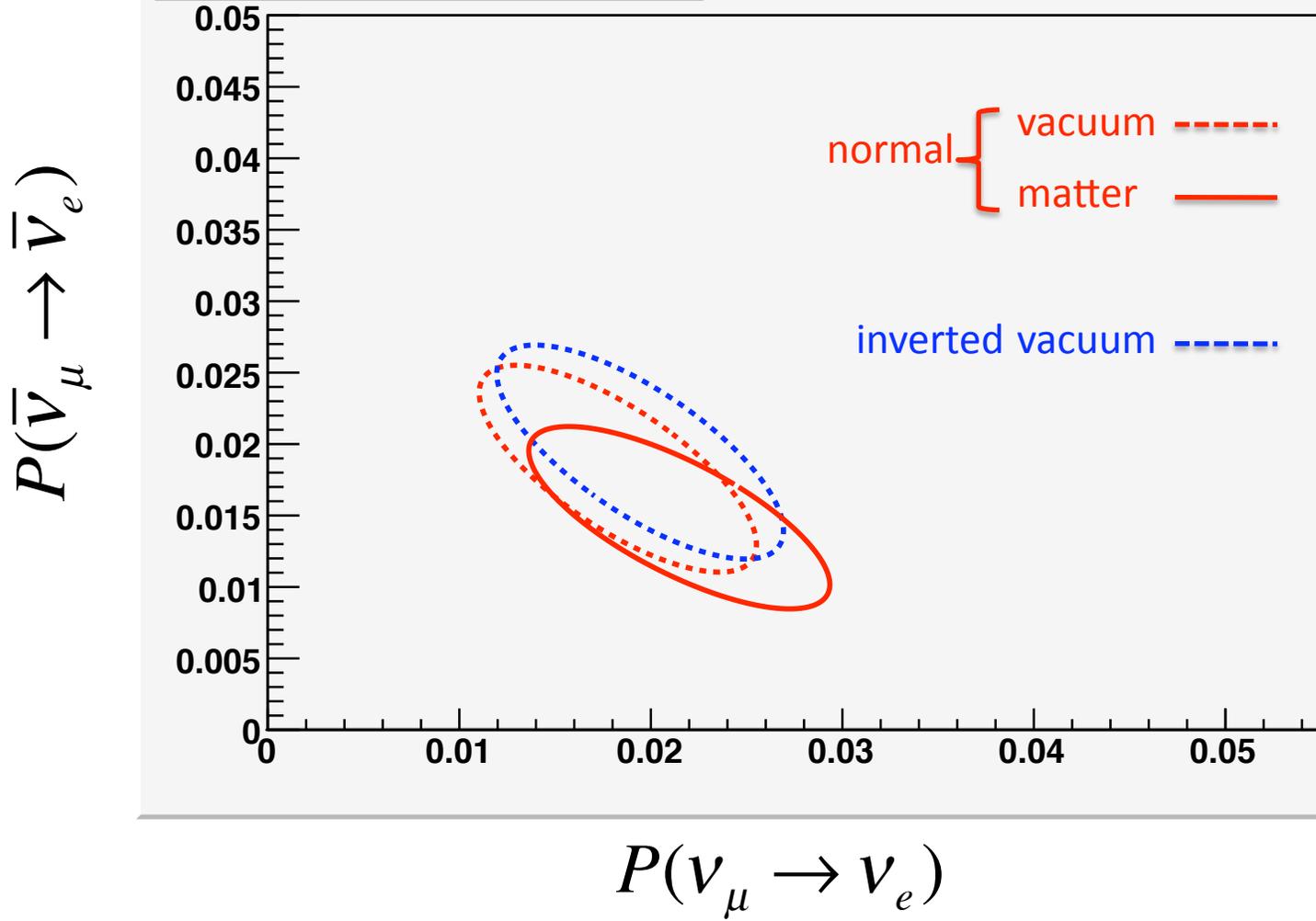
# Neutrino vs. anti-neutrino bi-probability plots



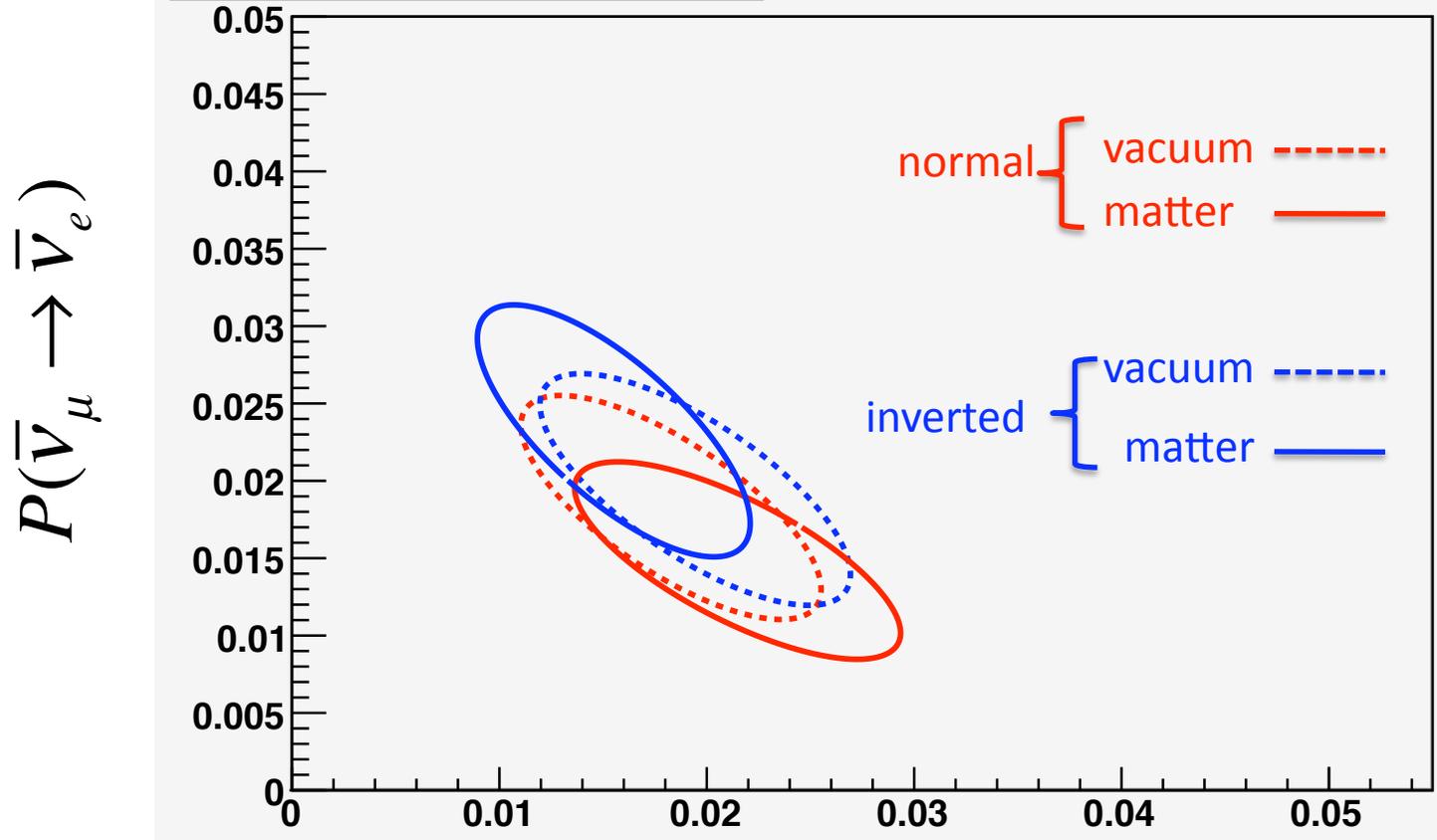
$\sin^2 2\theta_{13} = 0.04$   $L = 810$  km



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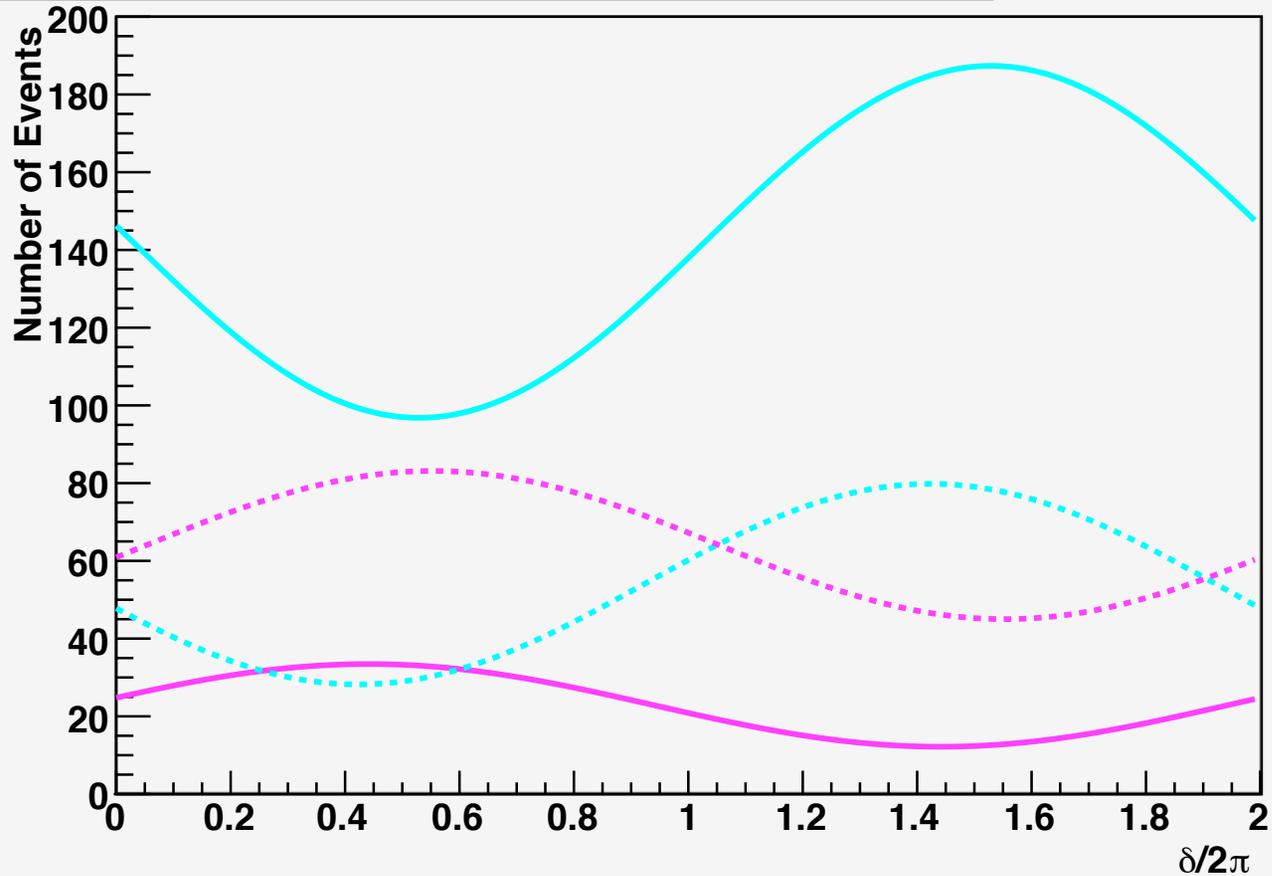
$\sin^2 2\theta_{13} = 0.04$   $L = 810$  km



$P(\nu_\mu \rightarrow \nu_e)$

# Why so large?

$\sin^2 2\theta_{13} = 0.04$   $18 \times 10^{20}$   $\nu$ ,  $18 \times 10^{20}$   $\bar{\nu}$ , 20kT LAr



$\sin^2 2\theta_{13} = 0.04$   $18 \times 10^{20} \nu$ ,  $18 \times 10^{20} \bar{\nu}$ , 20kT LAr

